# Some Congruence Theorems for Closed Hypersurfaces in Riemann Spaces (Part III: Method baser on Voss' Proof) 

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## Some Congruence Theorems for Closed Hypersurfaces

## in Riemann Spaces

(Part III: Method based on Voss' Proof)

by Heinz Hopf $\dagger$ (Zürich) and Yoshie Katsurada (Sapporo)

## Introduction

An idea that gives congruence of two hypersurfaces concerning a transformation group by a relation between the invariant of the corresponding points of these hyper surfaces was first introduced by H. Hopf and K. Voss [1], that is, in that paper congruence relations of two closed curves on a plane and of two closed surfaces in 3-dimensional euclidean space have been given by the relation of the mean curvatures.
K. Voss has generalized these theorems to hypersurfaces in an $(m+1)$-dimensional euclidean space $(m+1 \geqq 3)$ and also given the congruence relations in case of Gauss curvatures or the $r$-th mean curvatures $H_{r} r=1,2, \ldots, m$ [2]. A. Aeppli has developed analogous statements for a central transformation group (a homothetic transformation group with the center 0) [3].

The present authors wished to generalize these theorems to Riemann spaces. In the previous papers [4], [5], we gave the generalized theorems relating to the first mean curvature.

The purpose of the present paper is to investigate a general theorem relating to the Gauss curvature or the $r$-th mean curvature, that is, to generalize to an orientable Riemann space $R^{m+1}$ with constant Riemann curvature the following theorems given by K. Voss:

THEOREM (K. Voss). Let $W^{m}$ and $W^{m}$ be two orientable closed hypersurfaces in an $(m+1)$-dimensional euclidean space and let $p$ and $\bar{p}$ be the corresponding points of these hypersurfaces, and let $K(p)$ and $R(p)$ be that Gauss curvatures at these points respectively. Assume that the second fundamental forms of $W^{m}$ and $W^{m}$ are positive definite. If all straight lines ( $p \bar{p}$ ) are parallel to one another and if $K(p)=\bar{K}(\bar{p})$ for all $p \in W^{m}$, then the hypersurface $\bar{W}^{m}$ is produced from $W^{m}$ by simple translation in the direction of $(p \bar{p}) .\left(W^{m}\right.$ and $W^{m}$ are therefore congruent mod the translation group $)$.

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points respectively, for some $r=1,2, \ldots, m$. Assume that the second fundamental forms of $W^{m}$ and $\bar{W}^{m}$ are positive definite. If all straight lines $(p \bar{p})$ are parallel to one another and if $H_{r}(p)=\bar{H}_{r}(\bar{p})$ for all $p \in W^{m}$, then the hypersurface $W^{m}$ is produced from $W^{m}$ by simple translation in the direction of $(p \bar{p}) .\left(W^{m}\right.$ and $W^{m}$ are congruent mod the translation group.)

## § 1. Generalized Theorems

We suppose an ( $m+1$ )-dimensional orientable Riemann space with constant curvature $S^{m+1}$ of class $C^{v}(v \geqq 3)$ which admits an infinitesimal isometric transformation

$$
\begin{equation*}
\hat{x}^{i}=x^{i}+\xi^{i}(x) \delta \tau \tag{1.1}
\end{equation*}
$$

(where $x^{i}$ are local coordinate in $S^{m+1}$ and $\xi^{i}$ are the components of a contravariant vector $\xi$ ). We assume that orbits of the transformations generated by $\xi$ cover $S^{m+1}$ simply and that $\xi$ is everywhere continuous and $\neq 0$. Let us choose a coordinate system such that the orbits of the transformations are new $x^{1}$-coordinate curves, that is, a coordinate system in which the vector $\xi^{i}$ has components $\xi^{i}=\delta_{1}^{i}$, where the symbol $\delta_{j}^{i}$ denotes Kronecker's delta; then (1.1) becomes as follows

$$
\begin{equation*}
\hat{x}^{i}=x^{i}+\delta_{1}^{i} \delta \tau \tag{1.2}
\end{equation*}
$$

and $S^{m+1}$ admits a one-parameter continuous group $G$ of transformations which are 1-1-mappings of $S^{m+1}$ onto itself and are given by the expression $\hat{x}^{i}=x^{i}+\delta_{1}^{i} \tau$ in the new special coordinate system ([6]).

Now we consider two orientable closed hypersurfaces $W^{m}$ and $W^{m}$ of class $C^{v}$ imbedded in $S^{m+1}$ which are given as follows

$$
\left.\begin{array}{l}
W^{m}: x^{i}=x^{i}\left(u^{\alpha}\right) \quad i=1, \ldots, m+1 \quad \alpha=1, \ldots, m  \tag{1.3}\\
\bar{W}^{m}: \bar{x}^{i}=\bar{x}^{i}\left(u^{\alpha}\right)+\delta_{1}^{i} \tau\left(u^{\alpha}\right)
\end{array}\right\}
$$

where $u^{\alpha}$ are local coordinates of $W^{m}$ and $\tau$ is a continuous function attached to each point of the hypersurface $W^{m}$. We shall henceforth confine ourselves to Latin indices running from 1 to $m+1$ and Greek indices from 1 to $m$, and to two hypersurfaces $W^{m}$ and $\bar{W}^{m}$ which do not contain a piece of a hypersurface covered by the orbits of the transformations, which is expressed by $f\left(x^{2}, \ldots, x^{m+1}\right)=0$.

Then we can take the family of the hypersurfaces

$$
W^{m}(t)=(1-t) W^{m}+t W^{m} \quad 0 \leqq t \leqq 1
$$

generated by $W^{m}$ and $W^{m}$ whose points correspond along the orbits of the transformations where $W^{m}$ and $W^{m}$ mean $W^{m}(0)$ and $W^{m}(1)$ respectively. Thus according to (1.3), $W^{m}(t)$ is given by the expression

$$
\begin{equation*}
W^{m}(t): x^{i}\left(u^{\alpha}, t\right)=(1-t) x^{i}\left(u^{\alpha}\right)+t \bar{x}^{i}\left(u^{\alpha}\right) \quad 0 \leqq t \leqq 1 \tag{1.4}
\end{equation*}
$$

and (1.4) may be rewritten as follows

$$
\begin{equation*}
W^{m}(t): x^{i}\left(u^{\alpha}, t\right)=x^{i}\left(u^{\alpha}\right)+\delta_{1}^{i} t \tau\left(u^{\alpha}\right) \quad 0 \leqq t \leqq 1 \tag{1.5}
\end{equation*}
$$

Let us denote the normal unit vector of $W^{m}(t)$ by $n^{i}(t)$ and its derivative with respect to $t$ by $n^{\prime i}(t)$. Then $g_{i j}$ being the metric tensor of $S^{m+1}$ and differentiating the following relations with respect to $t$,

$$
g_{i j} n^{i}(t) \frac{\partial x^{j}(u, t)}{\partial u^{\alpha}}=0, \quad g_{i j} n^{i}(t) n^{j}(t)=1
$$

since the transformation group $G$ is isometric, that is, $\partial g_{i j} / \partial x^{1}=0$, we have

$$
\begin{align*}
& g_{i j} n^{\prime i} \frac{\partial x^{j}(u, t)}{\partial u^{\alpha}}+g_{i j} n^{i}(t) \frac{d}{d t}\left(\frac{\partial x^{j}(u, t)}{\partial u^{\alpha}}\right)=0,  \tag{1.6}\\
& g_{i j} n^{i}(t) n^{\prime j}(t)=0 \tag{1.7}
\end{align*}
$$

From (1.6), (1.7) and

$$
\frac{d}{d t}\left(\frac{\partial x^{i}(u, t)}{\partial u^{\alpha}}\right)=\delta_{1}^{i} \frac{\partial \tau}{\partial u^{\alpha}}
$$

we get

$$
\begin{equation*}
n^{\prime i}(t)=-g^{\alpha \beta}(t) \tau_{\alpha} \delta_{1}^{l} n_{l}(t) \frac{\partial x^{i}(u, t)}{\partial u^{\beta}} \tag{1.8}
\end{equation*}
$$

where $g^{\alpha \beta}(t)$ is the contravariant metric tensor of $W^{m}(t)$ and $\tau_{\alpha}$ means $\partial \tau / \partial u^{\alpha}$. Throughout this paper repeated lower case Latin indices call for summation 1 to $m+1$ and repeated lower case Greek indices for summation 1 to $m$. And also for its covariant differential along $W^{m}(t)$ we have

$$
\begin{equation*}
\delta n^{\prime i}(t)=d n^{\prime i}(t)+\Gamma_{j l}^{i} n^{\prime j}(t) x_{\gamma}^{l} d u^{\gamma} \tag{1.9}
\end{equation*}
$$

where $\Gamma_{j l}^{i}$ is the Christoffel symbol with respect to the metric tensor $g_{i j}$ of $S^{m+1}$ and $x_{\gamma}^{l}$ means $\partial x^{l}(u, t) / \partial u^{\gamma}$.

Let us give henceforth the derivative with respect to $t$ by the dash. Calculating $\left(\delta n^{i}\right)^{\prime}$, we have

$$
\begin{aligned}
\delta n^{i}=d n^{i}+ & \Gamma_{j l}^{i} n^{j}(t) x_{\gamma}^{l} d u^{\gamma} \\
\left(\delta n^{i}\right)^{\prime}=\left(d n^{i}\right)^{\prime} & +\left(\Gamma_{j l}^{i}\right)^{\prime} n^{j}(t) x_{\gamma}^{l} d u^{\gamma} \\
& +\Gamma_{j l}^{i} n^{\prime j}(t) x_{\gamma}^{l} d u^{\gamma}+\Gamma_{j l}^{i} n^{j}(t)\left(x_{\gamma}^{l}\right)^{\prime} d u^{\gamma}
\end{aligned}
$$

since $G$ is isometric, that is, $\partial g_{i j} / \partial x^{1}=0$, we have $\partial \Gamma_{l j}^{i} / \partial x^{1}=0$. Consequently we obtain the following relation between $\delta n^{\prime i}$ and $\left(\delta n^{i}\right)^{\prime}$

$$
\begin{equation*}
\left(\delta n^{i}\right)^{\prime}=\delta n^{i}+\Gamma_{j 1}^{i} n^{j}(t) \tau_{\gamma} d u^{\gamma} \tag{1.10}
\end{equation*}
$$

We claim that the following theorems hold

THEOREM 1.1. Let $K$ and $\bar{K}$ be the Gauss curvature of $W^{m}$ and $W^{m}$ respectively. Assume that the second fundamental form of $W^{m}(t), 0 \leqq t \leqq 1$ is positive definite. If the relation $K=\bar{K}$ holds for each point $p \in W^{m}$, then $W^{m}$ and $\bar{W}^{m}$ are congruent $\bmod G$.

Proof. We consider the following differential form of degree $m-1$ attached to each point $p$ on the hypersurface $W^{m}(t)$

$$
\left.\begin{array}{rl}
\left(\left(n^{\prime}, \delta_{1} \tau, \delta n, \ldots, \delta n\right)\right) \stackrel{\text { def. }}{\equiv} & \sqrt{g}\left(n^{\prime}, \delta_{1} \tau, \delta n, \ldots, \delta n\right)  \tag{1.11}\\
= & (-1)^{m-1} \sqrt{ } g\left(n^{\prime}, \delta_{1} \tau, x_{\alpha_{1}}, \ldots, x_{\alpha_{m-1}}\right) \\
& \times b_{\beta_{1}}^{\alpha_{1}}(t) \ldots b_{\beta_{m-1}}^{\alpha_{m-1}}(t) d u^{\beta_{1}} \wedge \ldots \wedge d u^{\beta_{m-1}}
\end{array}\right\}
$$

where $g$ is the determinant of the metric tensor $g_{i j}$ of $S^{m+1}$, the symbol () means a determinant of order $m+1$ whose columns are the components of respective vectors, $b_{\alpha \beta}(t)$ is the second fundamental tensor of $W^{m}(t)$ and $b_{\alpha}^{\beta}(t)$ denotes $b_{\alpha \gamma}(t) g^{\beta \gamma}(t)$.

Then the exterior differential of the differential form (1.11) becomes as follows

$$
\begin{align*}
& d\left(\left(n^{\prime}, \delta_{1} \tau, \delta n, \ldots, \delta n\right)\right)=\left(\left(\delta n^{\prime}, \delta_{1} \tau, \delta n, \ldots, \delta n\right)\right)  \tag{1.12}\\
& \quad+\left(\left(n^{\prime}, \delta\left(\delta_{1}\right) \tau, \delta n, \ldots, \delta n\right)\right)+\left(\left(n^{\prime}, \delta_{1} d \tau, \delta n, \ldots, \delta n\right)\right)
\end{align*}
$$

because since $S^{m+1}$ is a space of constant curvature, we have

$$
\left(\left(n^{\prime}, \delta_{1} \tau, \delta \delta n, \delta n, \ldots, \delta n\right)\right)=0
$$

Because $G$ is isometric, the quantity $n_{i}(t) \delta_{1}^{i} \sqrt{g^{*}(t)}$ is independent of $t$, where $g^{*}(t)$ means the determinant of $g_{\alpha \beta}(t)$, we have

$$
\begin{equation*}
\left(\left((\delta n)^{\prime}, \delta_{1} \tau, \delta n, \ldots, \delta n\right)\right)=(-1)^{m}(m-1)!K^{\prime} n_{i}(t) \delta_{1}^{i} \tau d A(t) \tag{1.13}
\end{equation*}
$$

where $d A(t)$ is the area element of $W^{m}(t)$, and using (1.8), we obtain

$$
\left.\begin{array}{l}
\left(\left(n^{\prime}, \delta_{1} d \tau, \delta n, \ldots, \delta n\right)\right)=(-1)^{m}(m-1)!  \tag{1.14}\\
\quad \times \frac{1}{\sqrt{g^{*}(t)}} B^{\alpha \beta}(t) \tau_{\alpha} \tau_{\beta}\left(n_{i}(t) \delta_{1}^{i}\right)^{2} \sqrt{\mathrm{~g}^{*}}(t) d A(t)
\end{array}\right\}
$$

where $B^{\alpha \beta}(t)$ means the cofactor of an element $b_{\beta \alpha}(t)$ in the determinant $\left|b_{\alpha \beta}(t)\right|$ divided by $g^{*}(t)$.

By making use of (1.10), (1.12), (1.13), (1.14) and the relation

$$
\delta\left(\delta_{1}^{i}\right)=\Gamma_{j 1}^{i} x_{\gamma}^{j} d u^{\gamma}
$$

we have

$$
\begin{aligned}
& d\left(\left(n^{\prime}, \delta_{1} \tau, \delta n, \ldots, \delta n\right)\right)=(-1)^{m}(m-1)!\left\{K^{\prime} n_{i}(t) \delta_{1}^{i} \tau d A(t)\right. \\
& \left.\quad+\frac{1}{\sqrt{g^{*}(t)}} B^{\alpha \beta}(t) \tau_{\alpha} \tau_{\beta}\left(n_{i}(t) \delta_{1}^{i}\right)^{2} \sqrt{g^{*}(t)} d A(t)\right\} \\
& \quad+\left(\left(n^{\prime}, \tau \Gamma_{j 1} x_{\gamma}^{j} d u^{\gamma}, \delta n, \ldots, \delta n\right)\right) \\
& \quad-\left(\left(\Gamma_{j 1} n^{j}(t) \tau_{\gamma} d u^{\gamma}, \delta_{1} \tau, \delta n, \ldots, \delta n\right)\right)
\end{aligned}
$$

Next we shall prove that

$$
\begin{equation*}
\left(\left(n^{\prime}, \tau \Gamma_{j 1} x_{\gamma}^{j} d u^{\gamma}, \delta n, \ldots, \delta n\right)\right)-\left(\left(\Gamma_{j 1} n^{j}(t) \tau_{\gamma} d u^{\gamma}, \delta_{1} \tau, \delta n, \ldots, \delta n\right)\right)=0 \tag{1.15}
\end{equation*}
$$

For the first term of the left-hand member of (1.15), making use of (1.8), we can see the following

$$
\begin{align*}
& \left(\left(n^{\prime}, \tau \Gamma_{j 1} x_{\gamma}^{j} d u^{\gamma}, \delta n, \ldots, \delta n\right)\right)=(-1)^{m-1} \tau n_{l}(t) \delta_{1}^{l} \\
& \quad \times\left(\left(\Gamma_{j 1} x_{\gamma}^{j}, g^{\alpha \beta}(t) \tau_{\beta} x_{\alpha}, x_{\alpha_{1}}, \ldots, x_{\alpha_{m-1}}\right)\right)  \tag{1.16}\\
& \quad \times b_{\beta_{1}}^{\alpha_{1}}(t) \ldots b_{\beta_{m-1}}^{\alpha_{m-1}}(t) d u^{\gamma} \wedge d u^{\beta_{1}} \wedge \ldots \wedge d u^{\beta_{m-1}}
\end{align*}
$$

Let $\varepsilon_{i_{1} \cdots i_{m+1}}$ and $\varepsilon_{\alpha_{1} \cdots \alpha_{m}}$ be the $\varepsilon$-symbol of $S^{m+1}$ and of $W^{m}(t)$ respectively, $\varepsilon_{i_{1} \ldots i_{m+1}} \stackrel{\text { def. }}{=} \sqrt{g} e_{i_{1} \ldots i_{m+1}}, \quad \varepsilon_{\alpha_{1} \ldots \alpha_{m}} \stackrel{\text { def. }}{=} \sqrt{g^{*}(t)} e_{\alpha_{1} \ldots \alpha_{m}}$,
the symbol $e_{i_{1} \cdots i_{m+1}}$ meaning plus one or minus one, depending on whether the indices $i_{1}, \ldots i_{m+1}$ denote an even permutation of $1,2, \ldots, m+1$ or odd permutation, and zero when at least any two indices have the same value, and also the symbol $e_{\alpha_{1} \cdots \alpha_{m}}$ meaning similarly for the indices $\alpha_{1}, \ldots, \alpha_{m}$ running from 1 to $m$.

Making use of the relation

$$
n_{i}(t) \varepsilon_{\alpha \alpha_{1} \ldots \alpha_{m-1}}=\varepsilon_{i i_{2} \ldots i_{m+1}} x_{\alpha}^{i_{2}} x_{\alpha_{1}}^{i_{3}} \cdots x_{\alpha_{m-1}}^{i_{m+1}}
$$

we have

$$
\begin{aligned}
& \left(\left(n^{\prime}, \tau \Gamma_{j 1} x_{\gamma}^{j} d u^{\gamma}, \delta n, \ldots, \delta n\right)\right)=(-1)^{m-1} \tau n_{l}(t) \delta_{1}^{l} \Gamma_{j 1}^{i} n_{i}(t) x_{\gamma}^{j} \tau_{\beta} g^{\beta d}(t) \\
& \quad \times \varepsilon_{\alpha \alpha_{1} \ldots \alpha_{m-1}} b_{\beta_{1}}^{\alpha_{1}}(t) \ldots b_{\beta_{m-1}}^{\alpha_{m-1}}(t) d u^{\nu} \wedge d u^{\beta_{1}} \wedge \ldots \wedge d u^{\beta_{m-1}} \\
& \quad=(-1)^{m-1} \tau n_{l}(t) \delta_{1}^{l} \Gamma_{j 1}^{i} n_{i}(t) x_{\gamma}^{j} \tau_{\beta} \varepsilon_{\alpha_{1} \ldots \alpha_{m-1}}^{\beta} \varepsilon^{\gamma \beta_{1} \ldots \beta_{m-1}} b_{\beta_{1}}^{\alpha_{1}}(t) \ldots b_{\beta_{m-1}}^{\alpha_{m-1}}(t) d A(t)
\end{aligned}
$$

and we can see easily the following relation

$$
\begin{aligned}
\varepsilon_{\alpha_{1} \ldots \alpha_{m-1}}^{\beta} \varepsilon^{\nu \beta_{1} \ldots \beta_{m-1}} b_{\beta_{1}}^{\alpha_{1}}(t) \ldots b_{\beta_{m-1}}^{\alpha_{m-1}}(t) & =\varepsilon^{\beta \gamma_{1} \ldots \gamma_{m-1}} \varepsilon^{\gamma \beta_{1} \ldots \beta_{m-1}} b_{\gamma_{1} \beta_{1}}(t) \ldots b_{\gamma_{m-1} \beta_{m-1}}(t) \\
& =(m-1)!B^{\beta \gamma}(t) .
\end{aligned}
$$

Since $B^{\beta \gamma}(t)$ is the symmetric tensor, we have

$$
\left.\begin{array}{l}
\left(\left(n^{\prime}, \tau \Gamma_{j 1} x_{\gamma}^{j} d u^{\gamma}, \delta n, \ldots \delta n\right)\right)=(-1)^{m-1}(m-1)!\tau n_{l}(t)  \tag{1.17}\\
\quad \times \delta_{1}^{l} \Gamma_{j i 1} n^{i}(t) x_{(\gamma}^{j} \tau_{\beta)} B^{\beta \gamma}(t) d A(t)
\end{array}\right\}
$$

where $\Gamma_{j i 1}$ means $g_{i l} \Gamma_{j 1}^{l}$ and the symbol $(\gamma \beta)$ denotes the symmetric part for the indices $\gamma$ and $\beta$.

On the other hand, we calculate the second term of the left-hand member of (1.15). Since $G$ is isometric, that is, $\partial g_{i j} / \partial x^{1}=0$, we have

$$
\begin{align*}
\Gamma_{j 1}^{l} n^{j}(t) n_{l}(t) & =\frac{1}{2} g^{l k}\left(\frac{\partial g_{k j}}{\partial x^{1}}+\frac{\partial g_{1 k}}{\partial x^{j}}-\frac{\partial g_{j 1}}{\partial x^{k}}\right) n^{j}(t) n_{l}(t)  \tag{1.18}\\
& =\frac{1}{2} \frac{\partial g_{k j}}{\partial x^{1}} n^{j}(t) n^{k}(t)=0,
\end{align*}
$$

and we can give the vector $\delta_{1}^{i}$ by the expression

$$
\begin{equation*}
\delta_{1}^{i}=n_{l}(t) \delta_{1}^{l} n^{i}(t)+\varphi^{\beta} x_{\beta}^{i} \tag{1.19}
\end{equation*}
$$

Substituting (1.19) in the second term of the left-hand member of (1.15) and making use of (1.18), we have

$$
\begin{align*}
& -\left(\left(\Gamma_{j 1} n^{j}(t) \tau_{\gamma} d u^{\gamma}, \delta_{1} \tau, \delta n, \ldots, \delta n\right)\right) \\
& \quad=-(-1)^{m-1} \tau n_{l}(t) \delta_{1}^{l}\left(\left(\Gamma_{j 1} n^{j}(t) \tau_{\gamma}, n, x_{\left.\left.\alpha_{1}, \ldots, x_{\alpha_{m-1}}\right)\right)} \quad \times b_{\beta_{1}}^{\alpha_{1}}(t) \ldots b_{\beta_{m-1}}^{\alpha_{m-1}}(t) d u^{\gamma} \wedge d u^{\beta_{1}} \wedge \ldots \wedge d u^{\beta_{m-1}} .\right.\right. \tag{1.20}
\end{align*}
$$

Let us take the relation

$$
\varepsilon_{\alpha \alpha_{1} \ldots \alpha_{m-1}} g^{\alpha \beta}(t) x_{\beta}^{j} g_{i j}=(-1)^{m} \varepsilon_{i i_{2} \ldots i_{m+1}} x_{\alpha_{1}}^{i_{2}} \ldots x_{\alpha_{m-1}}^{i_{m}} n^{i_{m+1}} .
$$

Then we have

$$
\left.\begin{array}{l}
-\left(\left(\Gamma_{j 1} n^{j}(t) \tau_{\gamma} d u^{\gamma}, \delta_{1} \tau, \delta n, \ldots, \delta n\right)\right)  \tag{1.21}\\
\quad=(-1)^{m-1}(m-1)!\tau n_{l}(t) \delta_{1}^{l} \Gamma_{i j 1} n^{i}(t) x_{(\beta}^{j} \tau_{\gamma)} B^{\beta \gamma}(t) d A(t)
\end{array}\right\}
$$

Thus from (1.17), (1.21) and $\Gamma_{i j 1}+\Gamma_{j i 1}=\partial g_{i j} / \partial x^{1}=0$, we can arrive at (1.15) as follows

$$
\begin{aligned}
& \left(\left(n^{\prime}, \tau \Gamma_{j 1} x_{\gamma}^{j} d u^{\gamma}, \delta n, \cdots, \delta n\right)\right)-\left(\left(\Gamma_{j 1} n^{j}(t) \tau_{\gamma} d u^{\gamma}, \delta_{1} \tau, \delta n, \ldots, \delta n\right)\right) \\
& \quad=(-1)^{m-1}(m-1)!\tau n_{l}(t) \delta_{1}^{l}\left(\Gamma_{i j 1}+\Gamma_{j i 1}\right) n^{i}(t) x_{(\gamma}^{j} \tau_{\beta)} B^{\beta \gamma}(t) d A(t)=0 .
\end{aligned}
$$

Finally we have

$$
\begin{align*}
& \frac{(-1)^{m}}{(m-1)!} d\left(\left(n^{\prime}, \delta_{1} \tau, \delta n, \ldots, \delta n\right)\right)=K^{\prime} n_{i}(t) \delta_{1}^{i} \tau d A(t) \\
& \quad+\frac{1}{\sqrt{g^{*}(t)}} B^{\alpha \beta}(t) \tau_{\alpha} \tau_{\beta}\left(n_{i}(t) \delta_{1}^{i}\right)^{2} \sqrt{ } g^{*}(t) d A(t) \tag{1.22}
\end{align*}
$$

Integrating both members of (1.22) over the interval $0 \leqq t \leqq 1$, we get

$$
\left.\begin{array}{l}
\frac{(-1)^{m}}{(m-1)!} d \int_{0}^{1}\left(\left(n^{\prime}, \delta_{1} \tau, \delta n, \ldots, \delta n\right)\right) d t=(R-K) n_{i}(0) \delta_{1}^{i} \tau d A(0)  \tag{1.23}\\
\quad+\sqrt{g^{*}(0)} \int_{0}^{1} g^{*}(t)^{-1 / 2} B^{\alpha \beta}(t) d t \tau_{\alpha} \tau_{\beta}\left(n_{i}(0) \delta_{1}^{i}\right)^{2} d A(0)
\end{array}\right\}
$$

Furthermore integrating both members of (1.23) over $W^{m}$ and applying Stokes' theorem, since $W^{m}$ is closed, we have

$$
\begin{aligned}
& \iint_{W^{m}}(\mathbb{R}-K) n_{i}(0) \delta_{1}^{i} \tau d A(0) \\
& \quad+\iint_{W^{m}} \sqrt{g^{*}(0)} \int_{0}^{1} g^{*}(t)^{-1 / 2} B^{\alpha \beta}(t) d t \tau_{\alpha} \tau_{\beta}\left(n_{i}(0) \delta_{1}^{i}\right)^{2} d A(0)=0
\end{aligned}
$$

making use of the hypothesis $K=K$, we obtain

$$
\iint_{W^{m}} \sqrt{g^{*}(0)} \int_{0}^{1} g^{*}(t)^{-1 / 2} B^{\alpha \beta}(t) d t \tau_{\alpha} \tau_{\beta}\left(n_{i}(0) \delta_{1}^{i}\right)^{2} d A(0)=0
$$

On the other hand, from that the second fundamental form of $W^{m}(t)$ is positive definite everywhere in $W^{m}(t), 0 \leqq t \leqq 1$, the quantity

$$
\sqrt{g^{*}(0)} \int_{0}^{1} g^{*}(t)^{-1 / 2} B^{\alpha \beta}(t) d t v_{\alpha} v_{\beta}
$$

becomes positive definite. From that two hypersurfaces $W^{m}$ and $\bar{W}^{m}$ do not contain a piece of a hypersurface covered by the orbits of transformations, a point on $W^{m}$ such that $n_{i}(0) \delta_{1}^{i}=0$ must be an isolate point. Moreover since $\tau$ is a continuous function of $W^{m}$, we have

$$
\tau=\text { constant }
$$

for all points of $W^{m}$. Consequently we can arrive at the following result

$$
W^{m} \equiv W^{m} \bmod G
$$

THEOREM 1.2. Let $H_{r}$ and $\bar{H}_{r}$ be the $r$-th mean curvature of $W^{m}$ and $\bar{W}^{m}$ respectively. Assume that the second fundamental form of $W^{m}(t), 0 \leqq t \leqq 1$, is positive definite. If the relation

$$
H_{r}=\bar{H}_{r}
$$

holds for each point $p \in W^{m}$, then $W^{m}$ and $W^{m}$ are congruent $\bmod G$.
Proof. We consider the following differential form of degree $m-1$ attached to each point $p$ on the hypersurface $W^{m}(t)$

$$
\left.\begin{array}{l}
((n^{\prime}, \delta_{1} \tau, \underbrace{\delta n, \ldots, \delta n}, d x, \ldots, d x)) \stackrel{\text { def. }}{\equiv} \sqrt{g}\left(n^{\prime}, \delta_{1} \tau, \delta n, \ldots, \delta n, d x, \ldots, d x\right)  \tag{1.24}\\
\quad=(-1)^{r-1} \sqrt{\bar{g}}\left(n^{\prime}, \delta_{1} \tau, x_{\alpha_{1}}, \ldots, x_{\alpha_{r-1}}, x_{\beta_{r}} \ldots x_{\beta_{m-1}}\right) \\
\quad \times b_{\beta_{r}}^{\alpha_{r}}(t) \ldots b_{\beta \sigma_{-1}}^{\alpha_{y}-1}(t) d u^{\beta_{1}} \wedge \ldots \wedge d u^{\beta_{r-1}} \wedge d u^{\beta_{r}} \wedge \ldots \wedge d u^{\beta_{m-1}}
\end{array}\right\}
$$

The exterior differential of the differential form (1.24) becomes as follows

$$
\begin{aligned}
& d\left(\left(n^{\prime}, \delta_{1} \tau, \delta n, \ldots, \delta n, d x, \ldots, d x\right)\right)=\left(\left(\delta n^{\prime}, \delta_{1} \tau, \delta n, \ldots, \delta n, d x, \ldots, d x\right)\right) \\
& \quad+\left(\left(n^{\prime}, \delta\left(\delta_{1}\right) \tau, \delta n, \ldots, \delta n, d x, \ldots, d x\right)\right)+\left(\left(n^{\prime}, \delta_{1} d \tau, \delta n, \ldots, \delta n, d x, \ldots, d x\right)\right)
\end{aligned}
$$

because since $S^{m+1}$ is a space of constant curvature, it follows that

$$
\left(\left(n^{\prime}, \delta_{1} \tau, \delta \delta n, \ldots, \delta n, d x, \ldots, d x\right)\right)=0
$$

and also we have

$$
\left(\left(n^{\prime}, \delta_{1} \tau, \delta n, \ldots, \delta n, \delta d x, \ldots, d x\right)\right)=0
$$

Making use of (1.8), we have

$$
\begin{align*}
& \left(\left(n^{\prime}, \delta_{1} d \tau, \delta n, \ldots, \delta n, d x, \ldots, d x\right)\right) \\
& \quad=(-1)^{r-1} g^{\alpha \beta}(t) \tau_{\alpha} n_{i}(t) \delta_{1}^{i}\left(\left(\delta_{1} \tau_{\gamma}, x_{\beta}, x_{\alpha_{1}}, \ldots, x_{\alpha_{r-1}}, x_{\alpha_{\alpha_{2}}}, \ldots, x_{\alpha_{m-1}}\right)\right) \\
& \quad \times b_{\beta_{1}}^{\alpha_{1}}(t) \ldots b_{\beta_{r-1}}^{\alpha_{r}-1}(t) d u^{\gamma} \wedge d u^{\beta_{1}} \wedge \ldots \wedge d u^{\beta_{r-1}} \wedge d u^{\alpha_{r}} \\
& \quad \wedge \ldots \wedge d u^{\alpha_{m-1}}  \tag{1.25}\\
& \quad=(-1)^{r-1} g^{\alpha \beta}(t) \varepsilon_{\beta \alpha_{1} \ldots \alpha_{r}-1 \alpha_{r} \ldots \alpha_{m-1}} \varepsilon^{\gamma \beta_{1} \ldots \beta_{r-1} \alpha_{r} \ldots \alpha_{m-1}} \\
& \quad \times b_{\beta_{1}}^{\alpha_{1}}(t) \ldots b_{\beta_{r-1}}^{\alpha_{r-1}}(t)\left(n_{i}(t) \delta_{1}^{i}\right)^{2} \tau_{\alpha} \tau_{\gamma} d A(t) .
\end{align*}
$$

On the other hand, from (1.10) we get

$$
\begin{align*}
& \left(\left(\delta n^{\prime}, \delta_{1} \tau, \delta n, \ldots, \delta n, d x, \ldots, d x\right)\right) \\
& \quad=\left(\left((\delta n)^{\prime}, \delta_{1} \tau, \delta n, \ldots, \delta n, d x, \ldots, d x\right)\right)  \tag{1.26}\\
& \quad-\left(\left(\Gamma_{j 1} n^{j}(t) \tau_{\gamma} d u^{\gamma}, \delta_{1} \tau, \delta n, \ldots, \delta n, d x, \ldots, d x\right)\right)
\end{align*}
$$

And after some calculations, we have

$$
\begin{equation*}
(-1)^{r} m!H_{r}^{\prime} n_{i}(t) \delta_{1}^{i} d A(t)=r\left(\left(\delta_{1},(\delta n)^{\prime}, \delta n, \ldots, \delta n, d x, \ldots, d x\right)\right) \tag{1.27}
\end{equation*}
$$

because $n_{i}(t) \delta_{1}^{i} d A(t)$ is independent of $t$ and $d x^{i}=\delta_{1}^{i} d \tau$, that is, the same direction to $\delta_{1}$. Moreover we can prove similarly the following relation as the proof of (1.15)

$$
\begin{align*}
& \left(\left(n^{\prime}, \delta\left(\delta_{1}\right) \tau, \delta n, \ldots, \delta n, d x, \ldots, d x\right)\right) \\
& \quad-\left(\left(\Gamma_{j 1} n^{j}(t) \tau_{\gamma} d u^{\gamma}, \delta_{1} \tau, \delta n, \ldots, \delta n, d x, \ldots, d x\right)\right)=0 \tag{1.28}
\end{align*}
$$

Then putting

$$
(m-1)!c_{(r)}^{\alpha \beta}=\varepsilon_{\alpha_{1} \ldots \alpha_{r-1} \alpha_{r} \ldots \alpha_{m-1}}^{\alpha} \varepsilon^{\beta \beta_{1} \ldots \beta_{r-1} \alpha_{r} \ldots \alpha_{m-1}} b_{\beta_{1}}^{\alpha_{1}}(t) \ldots b_{\beta_{r-1}}^{\alpha_{r}-1}(t)
$$

and using (1.25), (1.26), (1.27) and (1.28), we have

$$
\begin{align*}
& d\left(\left(n^{\prime}, \delta_{1} \tau, \delta n, \ldots, \delta n, d x, \ldots, d x\right)\right) \\
& \quad=\frac{(-1)^{r-1}}{r} m!H_{r}^{\prime} n_{i}(t) \delta_{1}^{i} \tau d A(t)  \tag{1.29}\\
& \quad+(-1)^{r-1}(m-1)!c_{(r)}^{\alpha \beta} \tau_{\alpha} \tau_{\beta}\left(n_{i}(t) \delta_{1}^{i}\right)^{2} d A(t)
\end{align*}
$$

Integrating both members of (1.29) over the interval $0 \leqq t \leqq 1$, and putting

$$
C_{(r)}^{\alpha \beta}=g^{*}(0)^{1 / 2} \int_{0}^{1} g^{*}(t)^{-1 / 2} c_{(r)}^{\alpha \beta} d t
$$

we have

$$
\begin{align*}
& m\left(\bar{H}_{r}-H_{r}\right) n_{i}(0) \delta_{1}^{i} \tau d A(0)+r C_{(r)}^{\alpha \beta} \tau_{\alpha} \tau_{\beta}\left(n_{i}(0) \delta_{1}^{i}\right)^{2} d A(0) \\
& \quad=\frac{r(-1)^{r-1}}{(m-1)!} d \int_{0}^{1}\left(\left(n^{\prime}, \delta_{1} \tau, \delta n, \ldots, \delta n, d x, \ldots, d x\right)\right) d t \tag{1.30}
\end{align*}
$$

Furthermore integrating both members of (1.30) over $W^{m}$ and applying Stokes' theorem

$$
\begin{aligned}
& \frac{m}{r} \iint_{W^{m}}\left(\bar{H}_{r}-H_{r}\right) n_{i}(0) \delta_{1}^{i} \tau d A(0)+\iint_{W^{m}}\left(n_{i}(0) \delta_{1}^{i}\right)^{2} C_{(r)}^{\alpha \beta} \tau_{\alpha} \tau_{\beta} d A(0) \\
& \quad=\frac{(-1)^{r-1}}{(m-1)!} \int_{\partial W^{m}} \int_{0}^{1}\left(\left(n^{\prime}, \delta_{1} \tau, \delta n, \ldots, \delta n, d x, \ldots, d x\right)\right) d t
\end{aligned}
$$

Since $W^{m}$ is closed, we have

$$
\frac{m}{r} \iint_{W^{m}}\left(\bar{H}_{r}-H_{r}\right) n_{i}(0) \delta_{1}^{i} \tau d A(0)+\iint_{W^{m}}\left(n_{i}(0) \delta_{1}^{i}\right)^{2} C_{(r)}^{\alpha \beta} \tau_{\alpha} \tau_{\beta} d A(0)=0
$$

using the hypotheses $H_{r}=\bar{H}_{r}$ and that the second fundamental form of $W^{m}(t), 0 \leqq t \leqq 1$, is positive definite, and from that two hypersurfaces $W^{m}$ and $\bar{W}^{m}$ do not contain a piece of a hypersurface covered by the orbits of transformations, we can arrive at

$$
\tau_{\alpha}=0
$$

for all points of $W^{m}$, consequently we have
$\tau=$ constant
for all points of $W^{m}$. Accordingly we can see the following result
$W^{m}=W^{m} \bmod G$.
This proof follows to the method of that due to K. Voss [2].
Remark. In an euclidean space, if $G$ is translation group, that is, a special isometric transformation group, Theorem 1.1 and Theorem 1.2 just coincide with theorems of K. Voss given in the introduction.

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