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## 9. THE RELATED EQUATION.

We are prepared now to make the construction toward which this entire discussion has been directed.

Consider the equation

$$L^*(u) = 0 . \quad (9.1)$$

with

$$L^*(u) = \frac{1}{T} \begin{bmatrix} m^*(\eta_1) & - & - & - & - & - & m^*(\eta_p) & m^*(u) \\ Dm^*(\eta_1) & - & - & - & - & - & Dm^*(\eta_p) & Dm^*(u) \\ - & - & - & - & - & - & - & - \\ - & - & - & - & - & - & - & - \\ D^{p-1}m^*(\eta_1) & - & - & - & - & - & D^{p-1}m^*(\eta_p) & D^{p-1}m^*(u) \\ l^*(m^*(\eta_1)) & - & - & - & - & - & l^*(m^*(\eta_p)) & l^*(m^*(u)) \end{bmatrix} . \quad (9.2)$$

$T$  being the determinant given in (8.4). This is clearly a differential equation of the  $n^{\text{th}}$  order in  $u$ , for which each one of the functions  $y_j(z, \lambda)$  and  $\eta_i(z, \lambda)$  is a solution. For if  $\eta_i$  is substituted for  $u$  two of the columns of the determinant (9.2) are the same, and if  $u$  is replaced  $y_j$  every element of the last column vanishes. Because the  $n$  solutions thus produced are linearly independent the solutions of the equation (9.1) are completely known.

The co-factor of the element  $l^*(m(u))$  in the formula (9.2) is the determinant  $T$ . The expansion of the formula thus gives it the aspect

$$L^*(u) = l^*(m^*(u)) - \sum_{v=1}^p \frac{T_v}{T} D^{p-v} m^*(u) , \quad (9.3)$$

where  $T_v$  is the determinant that is obtainable from the formula (8.4) by replacing its elements  $D^{p-v} m^*(\eta_j)$  by  $l^*(m^*(\eta_j))$ .

From the formula (8.5) it is seen that

$$l^*(m^*(\eta_j)) = \lambda^n \sum_{v=1}^p \frac{\tau_v(z, \lambda)}{\lambda^r} \cdot \frac{D^{u-1} v_j}{\lambda^{u-1}} \quad (9.4)$$

with

$$\tau_v(z, \lambda) = \sum_{k=0}^p \bar{\beta}_k(z, \lambda) \sigma_{v, r}^{(p-k)}(z, \lambda) . \quad (9.5)$$

The replacements which change  $T$  to  $T_v$  are thus seen to be ones which replace

$$\lambda^{n-v} \left\{ \delta_{p-v, j} + \frac{\sigma_{j, r}^{(p-v)}}{\lambda^r} \right\} \text{ by } \lambda^n \frac{\tau_v}{\lambda^r}.$$

It follows that

$$\frac{T_v}{T} = \lambda^v \frac{\theta_v(z, \lambda)}{\lambda^r},$$

with some function  $\theta_v(z, \lambda)$  which is bounded over the  $z$  and  $\lambda$  domains. This gives to the relation (9.3) the form

$$L^*(u) = l^*(m^*(u)) - \frac{1}{\lambda^r} \sum_{v=1}^p \lambda^v \theta_v D^{p-v} m^*(u). \quad (9.7)$$

With the substitution of the expression for  $D^{p-v} m^*(u)$ , as it may be obtained from (4.3) by writing  $\bar{\gamma}_{i-s}$  in the place of  $\gamma_{i-s}$ , it is found that

$$L^*(u) = l^*(m^*(u)) - \frac{1}{\lambda^r} \sum_{j=1}^n \lambda^j \omega_j(z, \lambda) D^{n-j} u, \quad (9.8)$$

with

$$\omega_j(z, \lambda) = \sum_{v=1}^p \sum_{s=0}^p \lambda^{-s} \binom{p-v}{s} \theta_v D^s \bar{\gamma}_{\mu-v-s}.$$

A comparison of this with the earlier result (6.6) shows that

$$L^*(u) = L(u) - \frac{1}{\lambda^r} \sum_{j=1}^n \lambda^j \{ \epsilon_j(z, \lambda) + \omega_j(z, \lambda) \} D^{n-j} u. \quad (9.9)$$

The equation (9.1), whose solutions are completely known, thus has coefficients which differ from those of the given equation (2.1) only by terms that are of at least the  $r^{\text{th}}$  degree in  $1/\lambda$ . It is, therefore, by definition, a related equation.

#### REFERENCES

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