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Starkey, John
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On the Relationship of Pericline and Albite Twinning to the Composition and Structural State of Plagioclase Feldspars

By John Starkey (London, Ont., Canada)*)

With 3 figures and 2 tables in the text

Abstract

Published crystallographic data from X-ray analyses of 105 plagioclase feldspars have been collected and, where possible, the complete lattice angles have been calculated. From these data the orientation of the rhombic section and the obliquity of albite twins have been determined and the dependence of these parameters on both composition and structural state is demonstrated.

It is shown how information concerning the history of a plagioclase crystal may be obtained by measuring the orientation of the composition planes of secondary pericline twins. The possible influence of the obliquity on twinning is also discussed.

It is suggested that the ease of development of glide twins in the plagioclase feldspars is related to the crystal structure.

INTRODUCTION

The secondary origin of polysynthetic albite and pericline twins in plagioclase feldspars was first recognised independently by JUDD (1885), LEHMAN (1884) and VAN WERVEKE (1883) and more recently by ALLING (1936) and EMMONS and GATES (1943) among others. The secondary origin of the twinning in chess-board albite has been demonstrated by STARKEY (1959). Polysynthetic twins have been produced experimentally in feldspars by FÖRSTNER (1884), Mügge and Heide (1931), Borg, HANDIN and Higgs (1959) and STARKEY and BROWN (1964). There is no

^{*)} Prof. J. Starkey, University of Western Ontario, Department of Geology, London, Ontario, Canada.

compelling evidence that polysynthetic albite and pericline twins in plagioclase feldspars can be of primary origin (i. e. growth twins). This view is opposed to that expressed by VANCE (1961).

The study of secondary twinning in feldspars is of interest not only for its own sake but also because of the information which can be gained about the history of the crystals and therefore of the host rock. Earlier studies along these lines have centered on the variations in the attitude of the twin elements relative to the crystal lattice as a function of both chemical composition and structural state. Unfortunately the earlier conclusions were based on very meagre data and it seems opportune at the present time to attempt to remedy this deficiency.

THE DATA

Published crystallographic data for a large number of plagioclase feldspars have been collected and, where possible, the direct and reciprocal lattice angles have been calculated. The results of this compilation are presented in Tables 1 and 2. The calculations were made with the I.B.M. 7040 Computer at the University of Western Ontario. The sources of the data recorded in Tables 1 and 2 are annotated below.

Table 1: Values for α , γ , α^* and γ^* have been taken from LAVES and CHAISSON (1950, Table 1); the data are all from albites. Specimen number 1¹) is a natural low albite and specimens 2 and 3 are the same material after heat treatment. Specimens 4 and 5 are synthetic.

Table 2: Reference 1²), GOLDSMITH and LAVES (1955, Table 1). Values of α^* , γ^* and β are given for a synthetic anorthite. The remaining lattice angles have been calculated from the following expressions:

$$\cos\beta^* = \left(\frac{\cos\alpha^*\cos\gamma^*}{\sin\gamma^*\sin\alpha^*} - \cos\beta\right)\sin\alpha^*\sin\gamma^*,\tag{1}$$

$$\cos\alpha = \frac{\cos\gamma^*\cos\beta^* - \cos\alpha^*}{\sin\gamma^*\sin\beta^*},\tag{2}$$

$$\cos\gamma = \frac{\cos\alpha^* \cos\beta^* - \cos\gamma^*}{\sin\alpha^* \sin\beta^*}.$$
(3)

Reference 2, BROWN (1960, Tables 2 and 3). BROWN measured α^* , γ^* and β for forty-five natural plagioclase feldspars selected from probable

¹) The specimen numbers used throughout are those listed in the first columns of Tables 1 and 2 under Js. No.

²) The reference numbers are listed in Table 2, column 2.

low temperature environments; he excluded specimens which showed peristerite unmixing. BROWN measured the same parameters for fifteen of the specimens after heating. The compositions of the feldspars were determined optically except for five specimens for which chemical analyses were available (BROWN, op. cit. Table 1), these are specimens 15, 46, 65, 76 and 87. Values for β^* , α and γ have been computed from the expressions given under Reference 1.

Reference 3, BROWN (1960, Table 4). BROWN measured α^* and γ^* and calculated β^* for two feldspars after successive heat treatments. The compositions were determined optically. Values for α and γ have been derived from expressions 2 and 3. Values for β have been calculated from the following expression:

$$\cos\beta = \frac{\cos\alpha^*\cos\gamma^* - \cos\beta^*}{\sin\alpha^*\sin\gamma^*}.$$
 (4)

Reference 4, COLE, SÖRUM and TAYLOR (1951, Table 4). Complete lattice parameters are given for three natural feldspars, an albite, a labradorite and an anorthite. The albite has been analysed chemically by SPENCER (1935—1937, Table 1, Specimen T); it occurs in micapegmatites and is probably of low temperature origin. The composition of the labradorite has been determined by EMMONS as $An_{56.1}$ from

Table 1

Data from LAVES and CHAISSON (1950), see text. In column 1 are listed the specimen numbers allocated in the present study; the specimen numbers in column 2 are from the original reference.

Js. No.	No.	$\mathbf{A}\mathbf{n}$	Type	Alpha	Gamma	Alpha Star	Gamma Star	Sigma	Phi
1	1	0.00	1	94.333	87.650	86.333	90.500	32.686	4.362
2	2	0.00	3	93.733	89.600	86.000	88.583	5.716	3.993
3	3	0.00	3	93.500	90.150	86.000	88.083	-2.149	3.990
4	5	0.00	4	93.433	90.033	86.133	88.250	-0.489	3.853
5	6	0.00	4	93.433	90.067	86.167	88.267	-1.002	3.845

Table 2

The references from which the data have been obtained are numbered in column 2; see text for details. The specimen numbers in column 1 have been assigned in order of increasing anorthite content. Column 3 lists the specimen numbers recorded in the original literature, a 0 indicates that no specimen number was given. Under AN are listed the compositions in Weight % Anorthite. Under Type are listed the nature of the specimen and its probable structural state, 1 = natural, probably ordered; 2 = natural, probably disordered; 3 = heat treated natural material, disordered; 4 = synthetic feldspar, disordered.

Phi	4.221 4.238 4.238 3.972 4.287	4.002 4.176 4.257 4.257	$\begin{array}{c} 4.358\\ 3.091\\ 0.000\\ 3.907\\ 3.907 \end{array}$	4.271 2.623 4.182 4.182 4.120	$\begin{array}{c} 3.974\\ 3.978\\ 3.978\\ 3.131\\ 3.513\\ 3.513\\ 3.395\\ 3.064\end{array}$	22.22.22.22.22.22.22.22.22.22.22.22.22.	3.967
Sigma	$\begin{array}{c} 32.381\\ 31.471\\ 33.296\\ 0.242\\ 32.944\end{array}$	-4.017 -2.637 -4.673 32.725 31.896	-29.832 -4.072 -5.126 -29.832 -4.158	$\begin{array}{c} 32.611\\ 32.108\\ -4.077\\ -3.394\\ 25.265\\ 26.402\end{array}$	$\begin{array}{c} 11.036\\ 1.877\\ 21.537\\ -5.524\\ -5.524\\ -4.828\\ -4.828\\ 515\\ -5.515\end{array}$	-5.899 -5.9985 -5.99855 -5.99855 -5.99455 -2.348855 -2.879 -2.879 -5.8055 -5.8299 -5.80555 -5.80555 -5.80555 -5.805555 -5.805555555555555555555	13.980
Gamma Star	$\begin{array}{c} 90.433\\ 90.367\\ 90.500\\ 88.250\\ 90.480\end{array}$	87.970 87.967 87.900 90.450 90.450	90.467 90.467 88.433 90.000 90.000 88.017	90.467 90.417 88.683 89.900 89.900 89.990	88.940 88.350 89.650 88.350 88.250 88.250 88.233 88.233 88.233 88.233 88.383	88,461 88,461 88,563 88,5660 88,5660 88,5563 80,556 80,5	89.150
Beta Star	63.586 63.567 63.433 63.450 63.450 63.510	63.500 63.481 63.530 63.567 63.567	63.633 63.660 63.667 63.883 63.833 63.833 63.633 63.633	63.740 63.522 63.522 63.560 63.560 63.422 63.500	$\begin{array}{c} 63.540\\ 63.510\\ 63.502\\ 63.704\\ 63.417\\ 63.410\\ 63.450\\ 63.450\\ 63.450\\ 63.650\end{array}$	63.583 63.633 63.633 63.612 63.612 63.612 63.582 63.547 63.547 63.570 63.580 63.580 63.520 63.500 6	63.670
Alpha Star	86.433 86.383 86.387 86.033 86.033 86.410	86.010 85.828 85.928 86.417 86.417	86.333 86.333 86.333 86.3517 87.5517 8	86.400 86.350 87.383 85.950 86.217 86.310	$\begin{array}{c} 86.100\\ 86.267\\ 86.267\\ 86.267\\ 86.883\\ 86.833\\ 86.510\\ 86.510\\ 86.517\\ 86.517\\ 86.517\\ 86.517\end{array}$	87.167 87.167 87.167 86.217 86.233 86.183 86.183 86.183 86.183 86.183 86.183 86.183 86.183 86.183 86.183 86.183 86.233 86.233	86.150
Gamma	87.741 87.789 87.617 89.983 87.670	90.280 90.192 87.700 87.700	87.650 90.219 90.220 90.000 90.000 90.220	87.700 87.712 90.186 90.240 88.216 88.216	$\begin{array}{c} 89.240\\ 89.240\\ 90.301\\ 90.339\\ 90.239\\ 90.299\\ 90.294\\ 90.294\end{array}$	90.301 90.345 90.345 90.2655 88.650 90.200 90.200 98.88 88.89 90.200 90.200 90.276 88.88 90.200 90.276 90.200 90.200 90.276 88.89 90.200 90.200 88.89 90.200 90.200 88.80 90.200 90.200 88.800 80.000 80.000 80.000 80.000 80.000 80.000 80.000 80.000 80.000 80.000 80.000 80.000 80.000 80.000 80.000 80.000 80.000 80.0000 80.0000 80.0000 80.0000 80.0000 80.0000 80.0000 80.00000 80.00000000	89.042
Beta	116.500 116.517 116.517 116.667 116.500	116.447 116.390 116.517 116.517	116.367 116.358 116.117 116.167 116.167	116.350 116.557 116.557 116.033 116.633 116.633	$\begin{array}{c} 116.450\\ 116.440\\ 116.533\\ 116.533\\ 116.533\\ 116.543\\ 116.543\\ 116.543\\ 116.543\\ 116.305\end{array}$	$\begin{array}{c} 116.375\\ 116.326\\ 116.326\\ 116.326\\ 116.333\\ 116.417\\ 116.450\\ 116.450\\ 116.283\\ 116.283\\ 116.333\\ 116.333\\ 116.333\\ 116.417\\ 116.4$	116.333
Alpha	94.199 94.223 94.317 93.567 94.360	93.450 93.649 93.470 94.233	94.339 94.333 92.099 90.000 93.367	94.246 94.287 92.269 93.520 94.181 94.120	93.830 93.620 93.620 93.144 93.144 92.997 92.997 92.903	992 992 992 992 992 992 992 992 992 992	93.875
Type		1 co +t +t -t -	- ന ന ന ന ന)	• ∞ ∞ ⊣ ∞ ∞ ∞ ∞ ∞		د
$\mathbf{A}\mathbf{n}$	0.00	0.000	0.50 0.50 0.50 0.50 0.50 0.50 0.50 0.50	22.00 22.00 11.00 11.00	111.20 111.20 111.20 111.20 113.00 113.00 113.00 113.00 113.00 113.00	$\begin{array}{c} 113.00\\ 132.00\\ 132.00\\ 142.00\\ 117.90\\$	22.00
N0.	27 46 10	-00803	0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0	$260 \\ 51 \\ 26 \\ 26 \\ 26 \\ 26 \\ 26 \\ 26 \\ 20 \\ 20$	111 888 111 111 111 111 111 111 111 111	222222222222222222222222222222222222222	1 67
Ref.	1 01 01 IO IO		N 20 m m m m x	00000000	0 0 0 0 0 0 0	0 8 8 9 9 9 9 9 9 9 7 9 9 9 9 9 9 9 9 9 9 9 9 9 9 9 9 9	4 64
J_{S} . No.	\$~%0¢		11116 11116 11116	2627737 26277	- 20 50 50 50 50 50 50 50 50 50 50 50 50 50	。	52 2

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Table 2

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Pericline and .	Albite Twinni	ng to the Co	mposition and Stru	uctural State	261
89570 88957 88957 88957 89570 805700 80570 805700 80570 80570 80570 80570 80570 80570 80570 8050	8.2.2.2.2.2.2.2.2.2.2.2.2.2.2.2.2.2.2.2	3.962 3.9140 3.9146 3.9148 3.9148 3.9589 3.95999 3.95999 3.95999 3.95999 3.95999 3.95999 3.95999 3.95999 3.95999 3.95999 3.95999 3.959999 3.959999 3.959999 3.959999 3.959999 3.9599999 3.959999999999		8283 844 844 844 845 845 855 855 855 855 855	4.293 4.293 4.308
$\begin{array}{c} -2.502\\ 15.010\\ -3.506\\ 12.341\\ 12.298\\ 10.651\\ 9.144\\ 8.235\\ -4.223\\ 11.145\end{array}$	-4.288 -4.298 -4.429 -3.183 -3.153 -5.028 -5.028 -5.156 -5.156 -5.156	6.953 5.953 5.520 5.556 5.520 5.220 5.220 5.220 5.220 5.220	-5.801	$\begin{array}{c} 1.625\\ -0.224\\ 3.215\\ 3.215\\ 3.215\\ -2.247\\ -4.018\\ -6.622\\ -7.018\\ -7.018\\ -9.397\\ -10.423\\ -10.423\end{array}$	-15.508 -16.478 -15.571
88,367 89,220 89,033 89,033 88,933 88,933 88,933 88,933 88,933 88,933 88,933 88,933 88,933 88,933 88,970	88.020 88.167 88.750 88.750 88.733 88.733 88.733 88.733 88.733 88.783 88.683 88.683	88.683 88.133 88.700 88.717 88.483 88.483 88.683 88.083 88.617	88,500 88,500 88,500 88,500 88,500 88,500 88,500 88,500 88,500 81,500 87,500 88,500 88,500 80,5000 80,5000 80,5000 80,5000 80,5000 80,5000 80,5000 80,50000000000	88.380 88.267 88.567 88.433 88.4117 88.117 88.117 88.117 88.117 88.117 88.1583 87.583 87.583 87.567 88.7563 88.7563 88.7563 88.7563 88.7563 88.7563 88.7563 88.7563 88.7563 88.7563 88.7563 88.75583 87.75583 87.755785 87.755787777777777	87.150 87.100 87.150
63.803 63.560 63.570 63.570 63.579 63.579 63.489 63.489 63.432 63.432 63.432 63.432 63.432 63.500	63.640 63.576 63.345 63.941 63.660 63.679 63.679 63.696 63.471	63.556 63.556 63.556 63.725 63.466 63.669 63.5669 63.5669 63.5669 63.5669	63.640 63.642 63.642 63.643 63.6440 63.6440 63.6440 63.6440 63.6440 63.6440 63.6440 63.6440 63.6440 63.6748 63.6748 63.6748	63.680 63.697 63.697 63.697 63.697 63.648 63.648 63.671 63.622 63.671 63.622 63.673 63.673 63.672 63.673 63.672 63.672	63.796 64.033 64.078
86.600 86.230 86.230 86.100 86.133 86.100 86.133 86.100 86.133 86.190 86.190 86.190 86.190	86.130 86.433 86.433 86.083 86.083 86.083 86.083 86.083 86.117 86.117 86.117	86.067 86.267 86.150 88.117 86.117 86.117 86.117 86.117 86.117 86.117	86.030 86.117 86.117 86.100 86.030 86.0000 86.000 86.000 86.000 86.000 86.000 86.000 86.000 86.000 86.000 86.000 86.000 86.00000 86.00000 86.0000000000	86.120 86.117 86.117 86.133 86.050 86.050 85.930 85.930 85.917 85.917 85.917 85.917	85.850 85.883 85.850
$\begin{array}{c} 90.149\\ 88.990\\ 89.147\\ 89.158\\ 89.158\\ 89.280\\ 89.373\\ 89.373\\ 89.373\\ 89.290\\ 89.250\\ \end{array}$	90.290 90.276 90.276 89.467 90.350 89.429 89.429 89.571	89.521 89.530 89.534 89.534 89.533 80.533 80	90,22,00,22,00,00,00,00,00,00,00,00,00,00	$\begin{array}{c} 89.890\\ 89.89.89.89.89.89.89.89.89.89.89.89.89.8$	91.146 91.217 91.155
116.150 116.450 116.450 116.333 116.333 116.500 116.550 116.550 116.550 116.250 116.440	116.300 116.367 116.633 116.633 116.633 116.417 116.283 116.283 116.333	116.333 116.383 116.417 116.483 116.483 116.483 116.267 116.283	116.290 116.300 116.300 116.300 116.300 116.300 116.383 115.300 116.383 115.250 116.383	116.250 116.250 116.250 116.250 116.250 116.250 116.250 116.253 116.283 116.283 116.283 116.283	116.083 115.850 115.800
$egin{array}{c} 92.986\\ 93.830\\ 93.873\\ 93.873\\ 93.752\\ 93.769\\ 93.750\\ 93.750\\ 93.750 \end{array}$	93.340 93.072 93.072 93.755 93.745 93.749 93.855 93.855	93.738 93.654 93.654 93.697 93.697 93.385 93.705 93.705	93.420 93.420 93.591 93.591 93.591 93.581 93.581 93.581 93.581 93.581	93.520 93.577 93.577 93.476 93.428 93.428 93.428 93.428 93.428 93.427 93.328 93.417 93.328 93.328	93.231 93.231 93.231
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0 1 ° ° ° ° ° ° ° ° ° ° ° ° ° ° ° ° ° °	1 1 4 4 7 8 8 8 8 8 8 8 8 8 8 8 8 8 8 8 8 8 8 8	888844% 0018664%	299893884444 298898898989	0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0	9 101
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chemical analysis and as An_{51} from the refractive index (see Cole et al., op. cit., Table 1). The composition of the anorthite has been given as approximately An_{100} on the basis of refractive index and extinction (GAY, 1953—1955, Table 1) and as An_{95-100} on the basis of universal stage measurements (KEMPSTER, MEGAW and RADASLOVICH, 1962, p. 1007). GAY (op. cit., p. 175) points out that the anorthite is of low temperature origin.

Reference 5, LAVES (1952, Table 2). The complete lattice parameters of both a low and high temperature albite are given.

Reference 6, DONNAY and DONNAY (1952, Tables 5 and 10). The complete lattice parameters are listed for a synthetic albite.

Reference 7, SMITH (1956, Table 4). SMITH has calculated complete lattice parameters from powder data for nine natural plagioclase feldspars (eight of low temperature origin and one of high temperature origin), for seven of these feldspars after heat treatment (one of them heated for two different periods of time) and for four synthetic plagioclase feldspars. The specimen numbers listed in column 3, Table 2 are consecutive numbers, 1 to 21, which have been allocated to SMITH's samples since his specimen numbers are unwieldy.

Reference 8, FERGUSON, TRAILL and TAYLOR (1958, Table 3). Complete lattice parameters are given for a natural low temperature albite and for a natural albite after heat treatment. The compositions of both materials were determined chemically by EMMONS (see FERGUSON et al., op. cit., Table 1).

PERICLINE TWINNING

In pericline twins the twin axis is the *b* crystallographic axis and the twin plane and composition plane are the rhombic section. The rhombic section is an irrational plane, it can be defined as the plane which includes both the *b* axis and the normal to *b* in (010). The orientation of the rhombic section is specified by the angle σ , which is the angle between the trace of the rhombic section and the trace of (001) on (010), σ is considered positive if the trace of the rhombic section on (010) lies between +a and +c. The angle can be calculated from the relationship $\cot \sigma = \cos \alpha^*/\cot \gamma$. Calculated values for σ are listed in Table 1, column 9 and Table 2, column 12. The values for σ are plotted against composition in Figure 1.

In Figure 1 the data points for albite-rich compositions are distributed in two groups, those representing ordered feldspars (i. e. of low temperature origin) lying above those for more disordered ones. The groups tend to merge towards more anorthite-rich compositions. Unfortunately there are no data for plagioclase feldspars of high temperature origin within the composition range An_{50} to An_{100} . However it is probable that the trend indicated by the more sodic feldspars continues until at An_{100} the

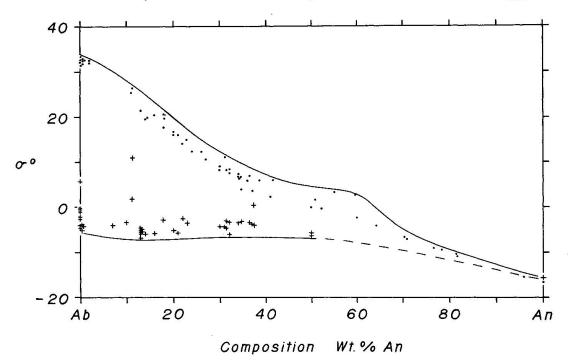


Fig. 1. The variation of σ with composition and structural state of plagioclase feldspars. The upper curve is for ordered feldspars and the lower one for disordered feldspars. The data are from Tables 1 and 2: dots represent ordered feldspars (Type 1) and crosses represent disordered feldspars (Types 2, 3 and 4). Brown's monalbites, specimens 19 and 20, are not included since the calculated value for σ has no meaning in a monoclinic crystal.

distinction between feldspars of high and low temperature origin disappears as is indicated by the data points for specimens 104 and 105 in Figure 1. This must be the case since anorthite occurs only in a highly ordered structural state. Since it is not certain that any of the feldspars considered here are in a state of maximum order or disorder, curves indicating the variation of σ with composition for ordered and disordered feldspars have been drawn along the "outer" margins of the data. These two curves are in general agreement with those of SMITH (1958, Figures 1 and 5), although the present curve for disordered feldspars indicates slightly larger negative values of σ for compositions more sodic than An₅₅, and both curves indicate slightly smaller negative values for

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anorthite-rich specimens. It should, however, be remembered that SMITH's curves are based on fewer data and, further, for compositions more calcic than An₅₅ his curves are actually those for Γ and not σ !

ALBITE TWINNING

In albite twinning the twin plane and composition plane are (010) and the twin axis is [010]. The angular misfit between the two individuals on either side of the twin plane is termed the obliquity and this can be defined as the angle between the normal to the twin plane and the lattice row quasi-normal to it. In albite twinning this is the angle between [010] and the *b* crystallographic axis, it is designated ϕ . The value of ϕ can be calculated from the expression, $\cos \phi = \sin \alpha^* \sin \gamma$. The computed values for ϕ are listed in Table 1, column 10 and Table 2, column 13.

In Figure 2 ϕ is plotted against composition, again curves for maximum ordered and maximum disordered feldspars have been drawn along

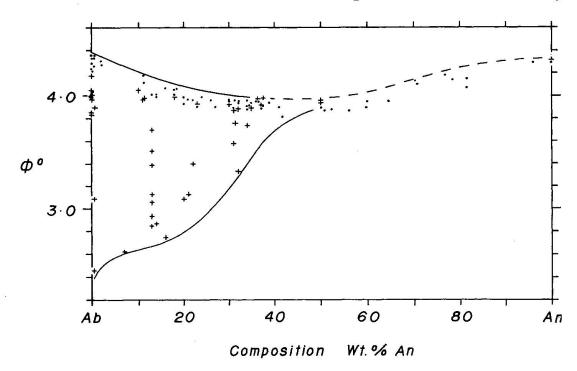


Fig. 2. The variation of ϕ with composition and structural state of plagioclase feldspars. The upper curve is for ordered feldspars and the lower one for disordered feldspars. The data are from Tables 1 and 2: dots represent ordered feldspars (Type 1) and crosses represent disordered feldspars (Types 2, 3 and 4). Brown's monalbites, specimens 19 and 20, are not included, they would plot at $\phi = 0^{\circ}$.

the "outer" margins of the data. Figure 2 shows that for ordered plagioclase feldspars ϕ is at a minimum around An₄₀ and increases towards An₀ and An₁₀₀. For the disordered feldspars the picture appears to be more complex, it is also, unfortunately, less complete. Data are lacking for compositions between An₅₀ and An₁₀₀ and it is therefore not possible to draw a curve with any certainty. However it is probable that the curves for ordered and disordered feldspars are sub-parallel and close together over this composition range, finally converging on anorthite (see the discussion under Pericline Twinning above). For disordered feldspars with compositions more sodic than An₅₀, ϕ decreases sharply and is 0° for BROWN's monalbite (specimens 19 and 20) as is dictated by the monoclinic symmetry.

The curves of Figure 2 differ considerably from those previously published (DONNAY 1940, GAY 1956—1958, SMITH 1958), the difference is particularly noticeable in the curve for disordered feldspars. In view of the well documented trend towards monoclinic symmetry exhibited by sodic plagioclase feldspars on heating it is thought that the trend of the curve for disordered feldspars proposed here is at least in the right direction.

CONCLUSION

The significance of the variation of σ and ϕ with composition and structural state have been discussed at length by SMITH (1958, 1962) so that detailed reiteration is not necessary here.

In the case of pericline twins the argument, briefly, is that, since the composition plane of the twins at the time of their formation was parallel to the rhombic section at that time, measurement of the composition plane and the determination of the present attitude of the rhombic section can give information on the thermal history of the feldspar and hence of the host rock.

In the case of secondary pericline twins which may result from deformation at different times during the cooling, and consequent ordering, of the plagioclase crystals, the values of σ for successively younger twins will decrease. The maximum value of σ will give an indication of the structural state in which the feldspar originally crystallized, remembering that it may have crystallized in a still more disordered state but that no twins developed at that time. The structural state in which the crystal originally grew can only be ascertained with certainty from the orientation of growth twins after the pericline law. Rocks undergoing heating and in which the plagioclase feldspars are becoming increasingly disordered might preserve evidence of this so that the earlier formed pericline twins would yield smaller values for σ . Further cooling could of course be accompanied by further generations of twins.

It has been thought that the ease of formation of albite twins reflects variations in the obliquity, ϕ (DONNAY 1940, GAY 1956—1958, SMITH 1958). The less the obliquity the easier the development of twins. In support of this idea SMITH (op. cit.) correlates his predictions with the observations of TURNER (1951) and GORAI (1951). However where the twinning results from deformation the problem is almost certainly more complex. Since glide twinning of an ordered plagioclase feldspar produces a markedly different structure in the twinned individual (STARKEY 1963), one can expect that where this structure is sufficiently unstable the crystal will resist twinning. On the basis of this Figure 3 has been prepared to show, in a purely qualitative way, how the ease of twinning can be expected to vary with composition and structural state. In deformation twinning the structural control will probably be more significant

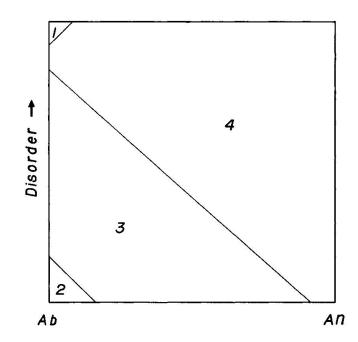


Fig. 3. The postulated effect of composition and structural state of plagioclase feldspars on the ease of formation of pericline and albite twins by gliding. Area 1 represents the possible field of monalbite, twinning is precluded by the monoclinic symmetry. Twinning is precluded on structural grounds in feldspars in area 2. In feldspars occurring in area 3 twinning is possible but not easy. Twinning is easy in feldspars occurring in area 4.

than variations in obliquity. The obliquity may exercise a control over the development of growth twins (see VANCE 1961).

Figure 3 is in general accord with the observations of TURNER (1951) and GORAI (1951). With increasing grade of metamorphism the plagioclase feldspars become more calcic and are in a structural state more favourable to twinning. Therefore frequency of twinning should increase with metamorphic grade. Secondary twinning should also be common on hornfelses and igneous rocks where temperatures are still higher and the structural state of the feldspars consequently more disordered.

Since the data which have been used to prepare Figures 1 and 2 were obtained at room temperature it might be supposed that the conclusions based on them apply only to feldspars which twin at room temperature. However SMITH (1958, p. 919) has pointed out that the orientation of the rhombic section is probably independent of temperature, this is perhaps also true in the case of the obliquity. As far as the boundaries indicated on Figure 3 are concerned, elevated temperatures and the accompanying increase in thermal vibration would tend to reduce the effect of the degree of structural order present so that the boundaries should shift towards the abscissus: the effects of this are not likely to be present in nature since prolonged heating would change the structural state. It is therefore felt that the conclusions presented above can be extended to natural environments.

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